



Study of the ablation of a solid wall by a liquid jet

ASTROFLU V

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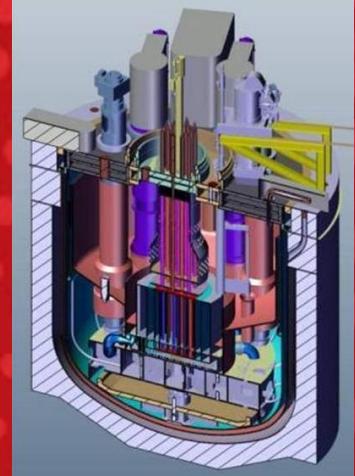
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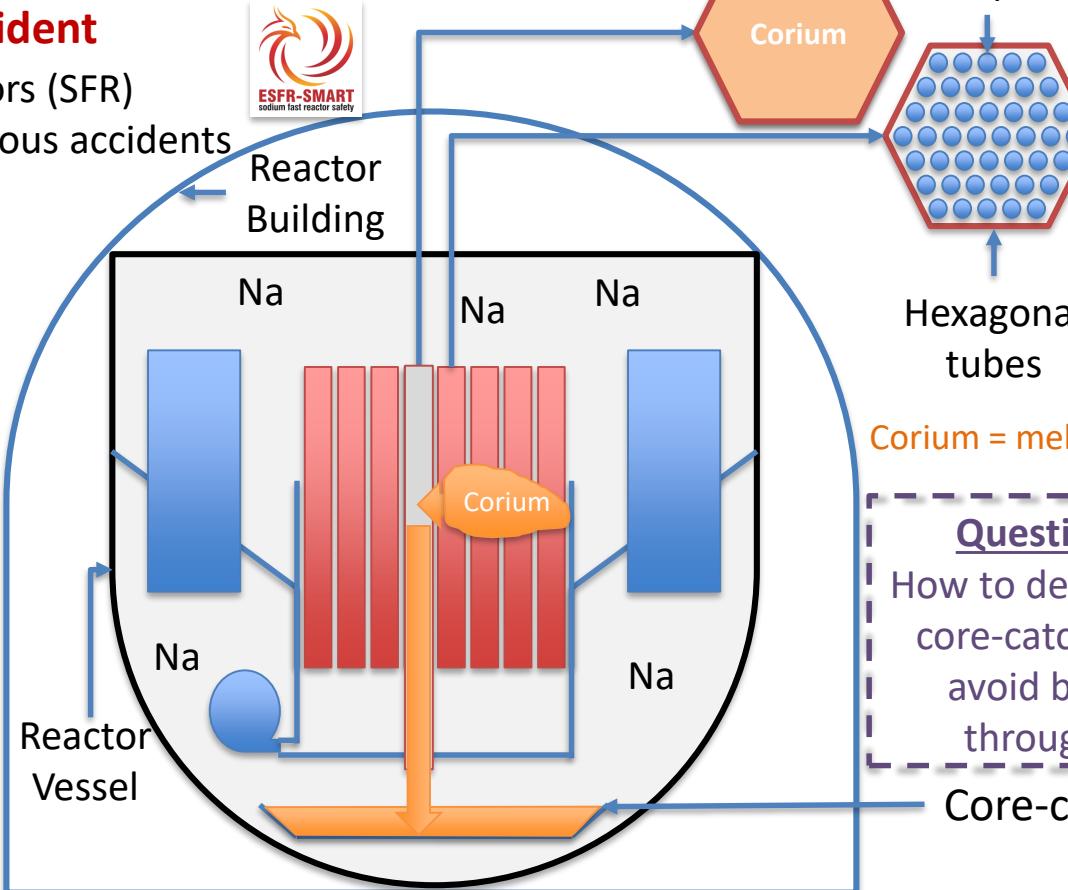


General context / Introduction

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Question :

How to design the core-catcher to avoid break through ?

Core-catcher

R&D context: Nuclear Severe Accident

- 4th generation: Sodium Fast Reactors (SFR)
- Integrate lessons learnt from previous accidents

Coping strategy for future SFR

Discharge tubes DT:

- Relocation toward Core-catcher as fast as possible

Core-catcher:

- Corium stabilization

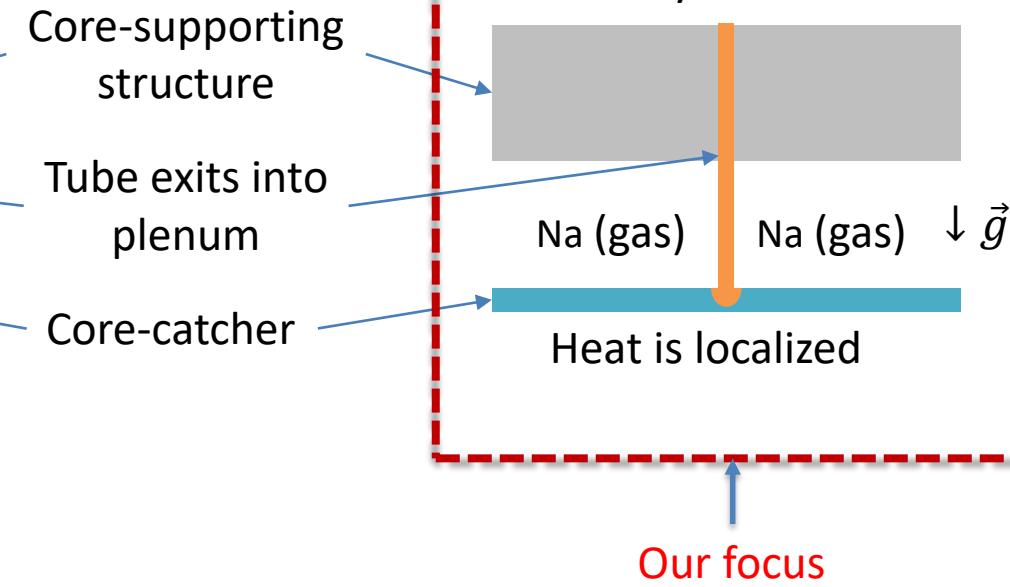
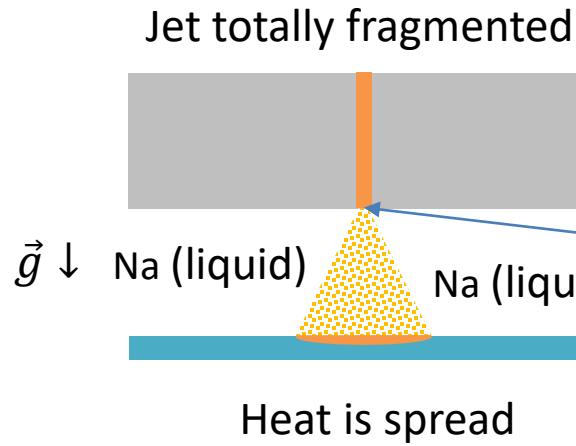
In-vessel retention

Advantages

- Reduce corium mass avoid recriticality
- Know corium path in advance



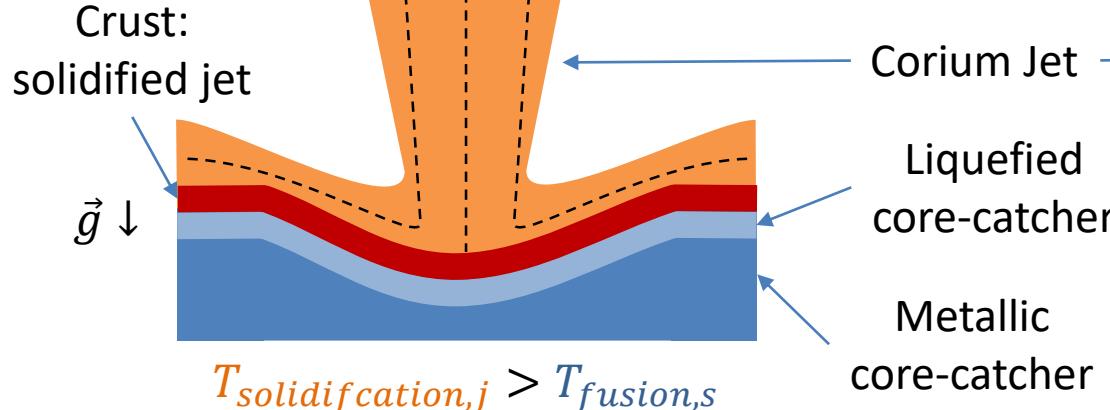
- **2 extreme cases at tubes outlet:**





- Coherent jet / 2 types of interactions:
 - Depending on jet / core-catcher relative compositions

Molten oxide jet

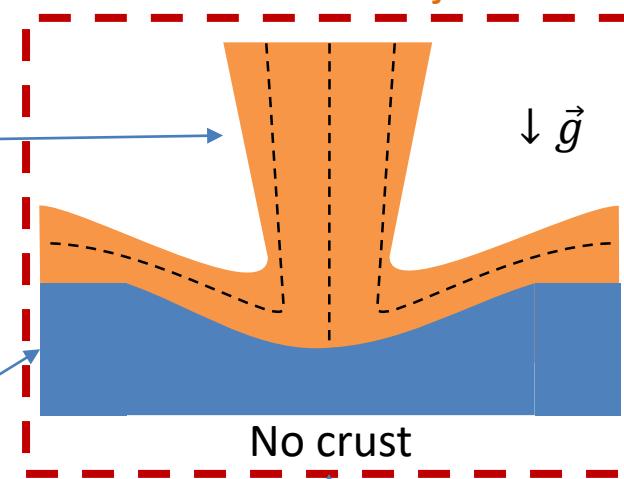


Crust acts as **thermal resistance**

Corium:

- oxide from nuclear fuel
- steel from structure

Molten metal jet

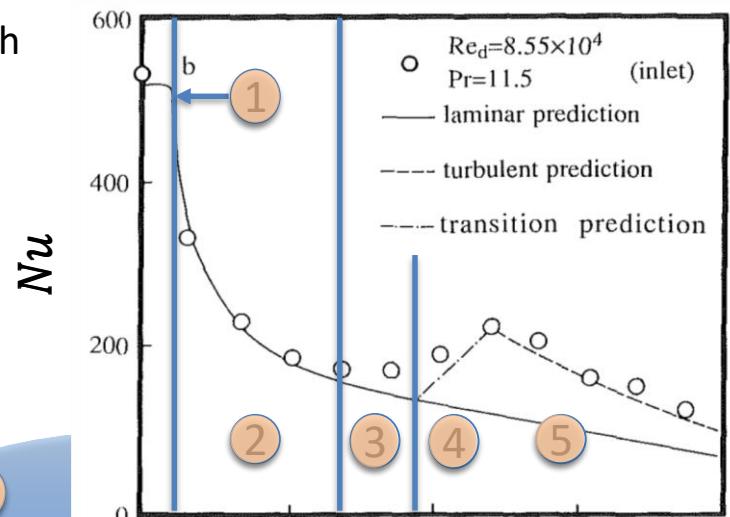
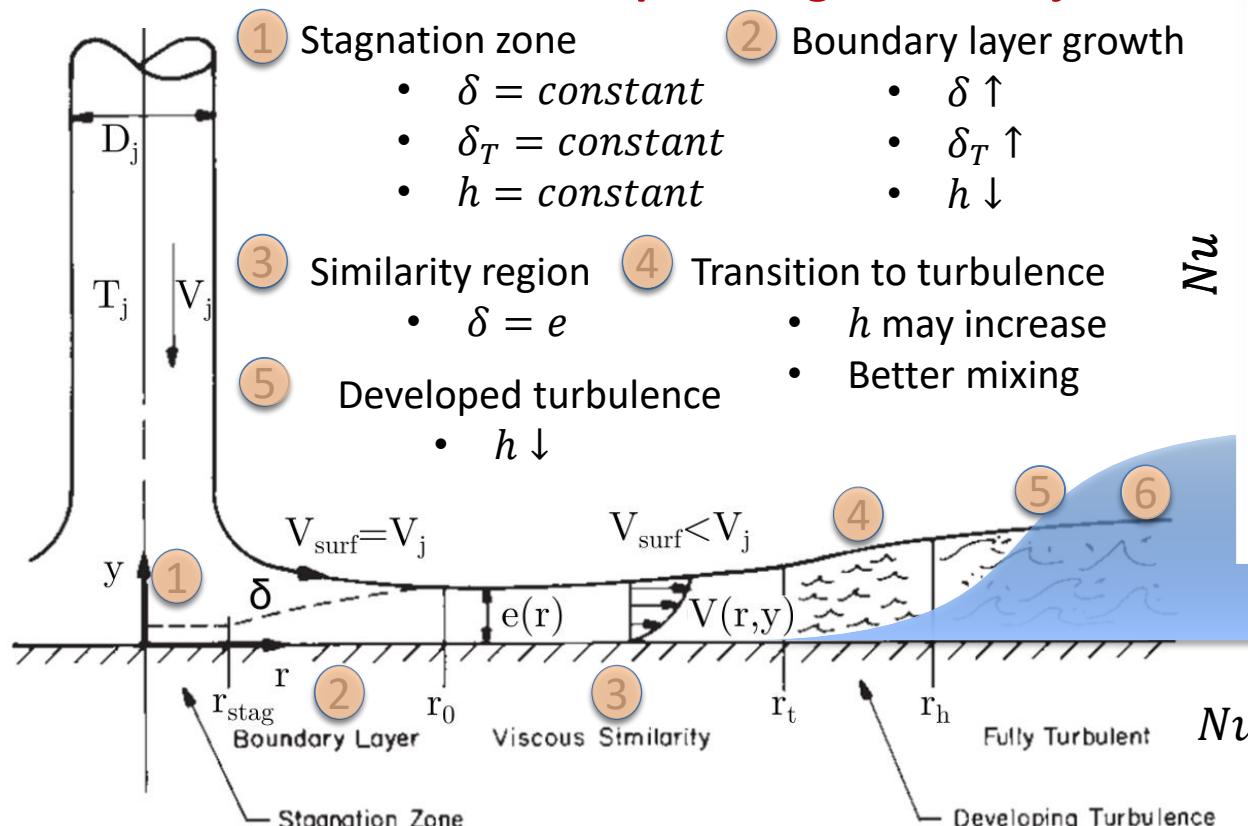


} Goals ← Focus

- Obtain quantitative data on ablation Jet / Solid same nature
- Improve understanding of ablation phenomenon



• Academic context: Jet spreading - laminar jets Watson (1964) + Lienhard (2006)



$$Nu = \frac{h D_j}{k_j}$$

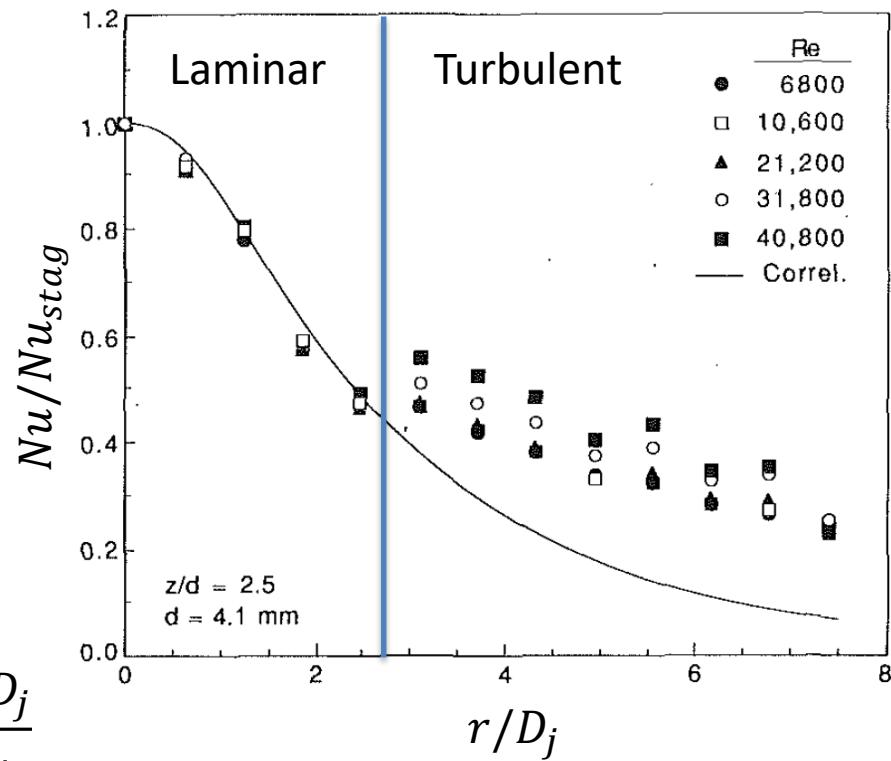
- $e \uparrow \frac{r}{D_j} \uparrow$
- $h \uparrow$ then \downarrow
- Can occur in any zone



- Academic context: heat transfer - turbulent jets Stevens et al. (1991)
+ Lienhard et al. (2006)

- SFR: Turbulent jets
- Flow laminarization at stagnation Tong (2003)
- Two regimes Stevens et al. (1991) + Lienhard (2006):
 - Laminar: Nu/Nu_{stag} independent of Re
 - Turbulent: Nu/Nu_{stag} depends on Re
- Transition:
 - Induced by jet surface oscillations
 - closer for turbulent jets
 - independent of Re
 - $h \uparrow$

$$Nu = \frac{h D_j}{k_j}$$





- **Dimensionless numbers**

Mechanical

Reynolds

Inertia vs. viscosity

$$Re = \frac{\rho_j D_j V_j}{\mu_j}$$

Froude

Inertia vs. gravity

$$Fr_j = \frac{V_j}{\sqrt{g D_j}}$$

Weber

Inertia vs. surface tension

$$We = \frac{\rho_j V_j^2 D_j}{\sigma_j}$$

Thermal

Nusselt

Convective vs. conductive transfer

$$Nu = \frac{h D_j}{k_j}$$

Melting number (Stefan)

Heat bring by jet vs. needed to melt

$$B = \frac{C_{p,j}(T_j - T_{s,f})}{L + C_{p,s}(T_{s,f} - T_{s,0})}$$

Stanton

Convective transfer vs. inlet heat flow

$$St = \frac{h}{\rho_j C_{p,j} V_j}$$

Prandtl

Compares mechanical / thermal BL growth

$$Pr = \frac{\mu_j C_{p,j}}{k_j}$$



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Experimental setup: HAnSoLO

Hot AblatioN of a SOLid by a Liquid - Observations

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• System studied: Water / Transparent ice

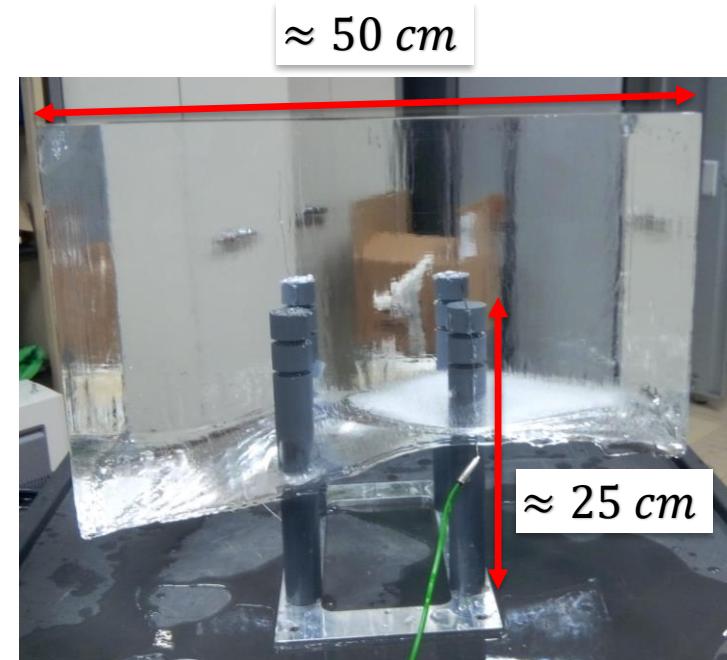
• Pros:

- Solid / Jet: same nature (no crust)
- Safe
- Cost effective
- Real-time visualizations

} Large number
of experiments

• Cons:

- Simulant not prototypical
- No undercooling (cracks)



• The experiment

Test referenced:
612

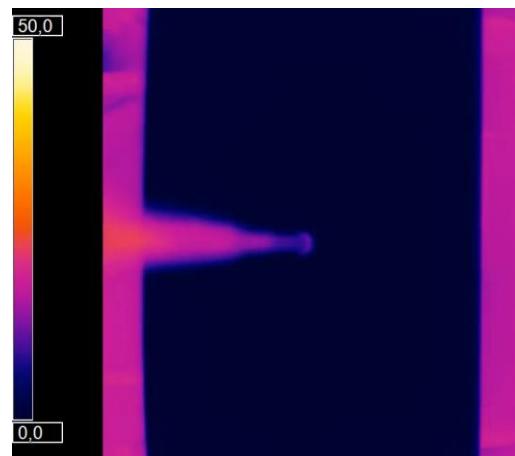
$$\begin{aligned}T_j &= 31 \text{ }^{\circ}\text{C} \\V_j &= 2.8 \text{ m/s} \\D_j &= 5.8 \text{ mm} \\Re &= 21\,000 \\Pr &= 5.30\end{aligned}$$





- Example of recordings

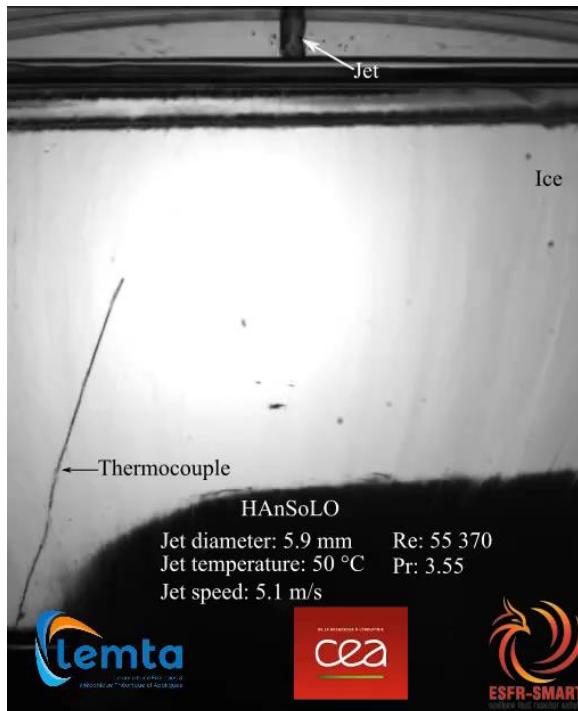
IR
camera



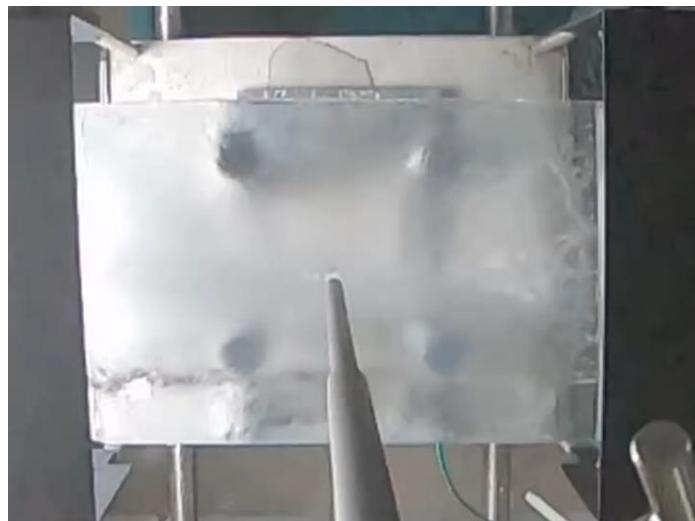
Test referenced: 623

$$\begin{aligned} T_j &= 50 \text{ }^{\circ}\text{C} \\ V_j &= 5.1 \text{ m/s} \quad Re = 55\,000 \\ D_j &= 5.9 \text{ mm} \quad Pr = 3.55 \end{aligned}$$

High speed
camera



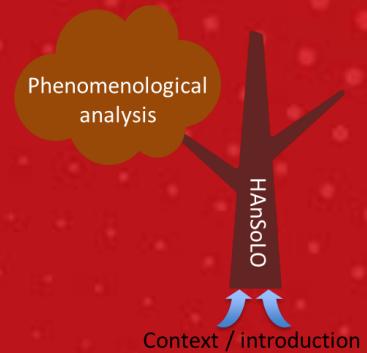
Standard
camera



POLYIMPERICEGENET



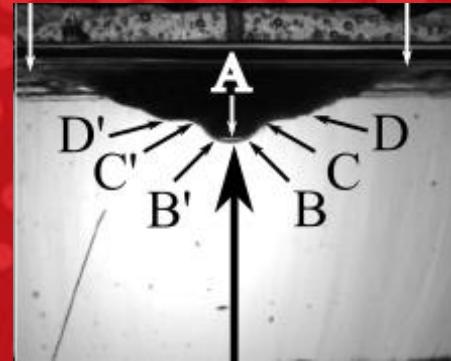
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Phenomenological analysis

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• Cavity actual shape

Sato et al. (1991)

steel/ste

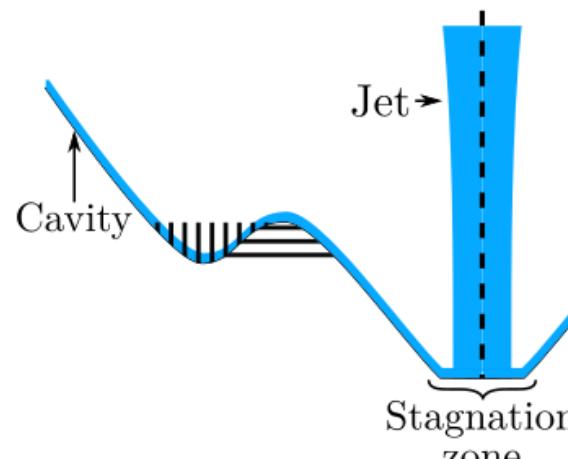
316L



$$D_j = 18.8 \text{ mm}$$

$$T_j = 1703 \text{ }^{\circ}\text{C}$$

$$V_j = 3.1 \text{ m/s}$$

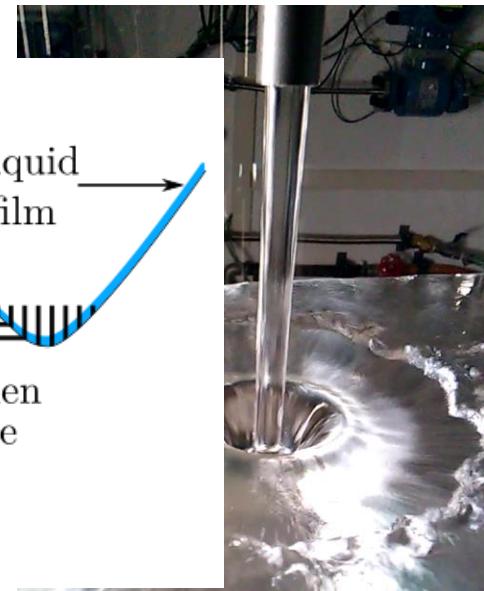


$$F$$

$$F$$

$$T_s$$

HAnSoLO
 Water / Ice



$$D_j = 8.1 \text{ mm}$$

$$T_j = 30 \text{ }^{\circ}\text{C}$$

$$V_j = 1.6 \text{ m/s}$$

$$Re = 19\,000$$

$$Pr = 5.38$$

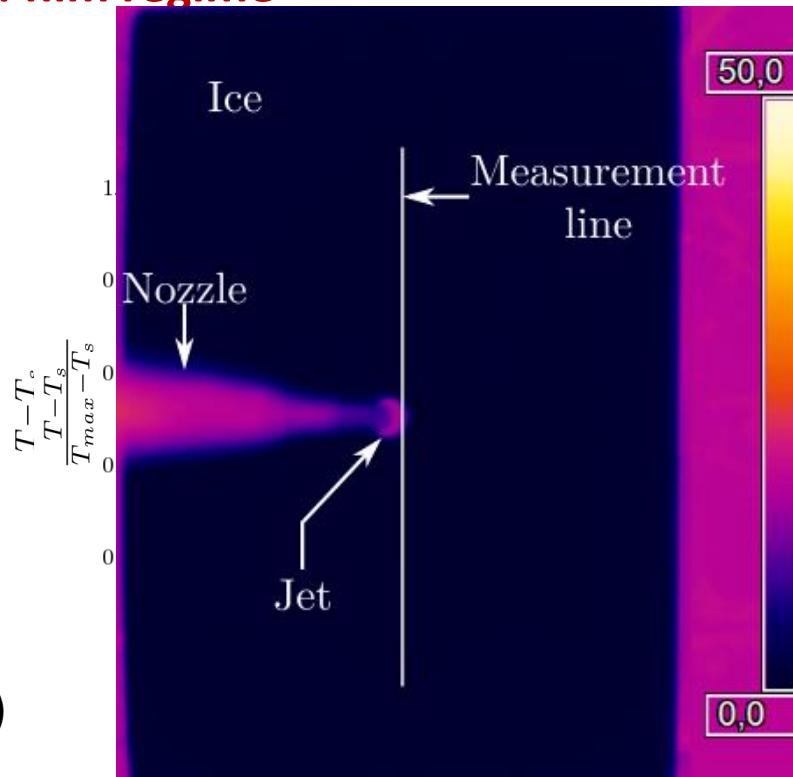
$$T_s = 0 \text{ }^{\circ}\text{C}$$

• Shape not due to considered system



• Film surface temperature: film regime

- Temperature along segment crossing jet
 - T independent of time
 - $T \downarrow$ at fixed r/D_j
 - For $V_j > 2.5 \text{ m.s}^{-1}$
 - Independent of Re
 - Due to turbulence
- Cf. Stevens et al. (1991)



623

$$T_j = 50 \text{ } ^\circ\text{C}$$

$$V_j = 5.1 \text{ m/s}$$

$$Re = 55\,000$$

$$Pr = 3.55$$

615

$$T_j = 30 \text{ } ^\circ\text{C}$$

$$V_j = 5.1 \text{ m/s}$$

$$Re = 77\,000$$

$$Pr = 5.35$$

634

$$T_j = 71 \text{ } ^\circ\text{C}$$

$$V_j = 7.6 \text{ m/s}$$

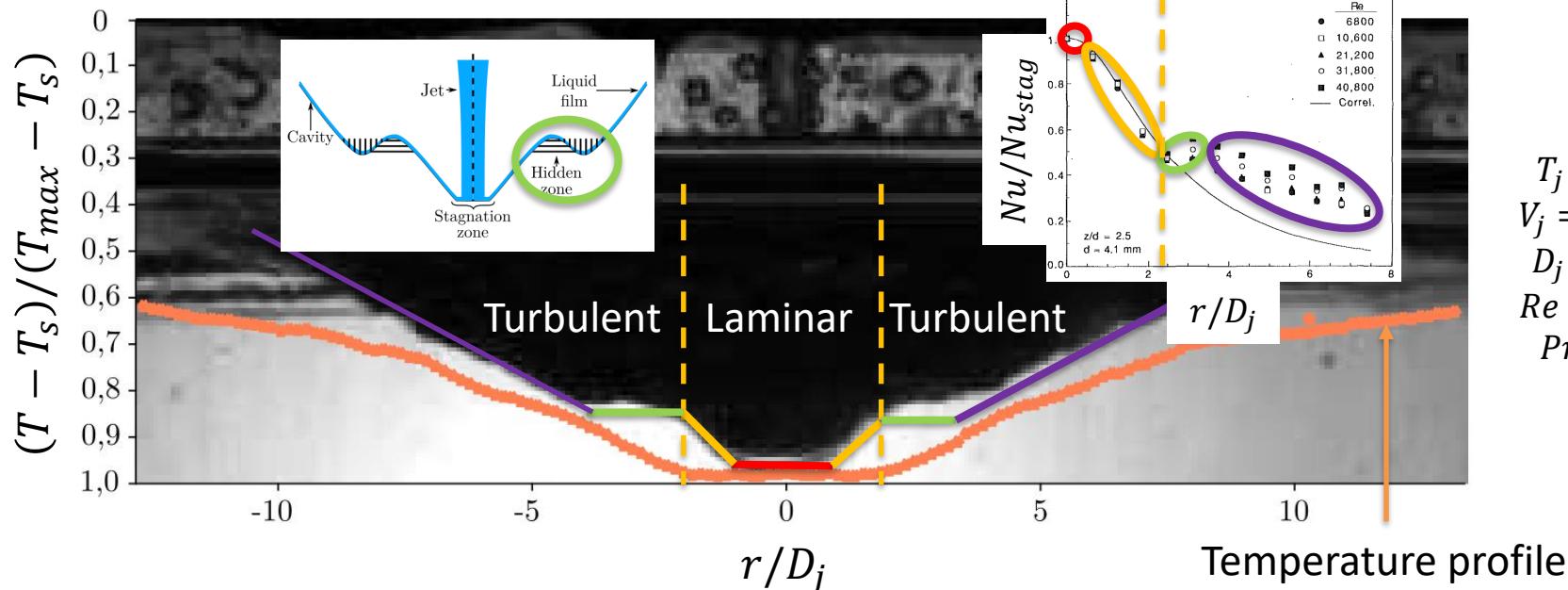
$$Re = 112\,000$$

$$Pr = 2.52$$



- Comparison: surface temperature / cavity shape

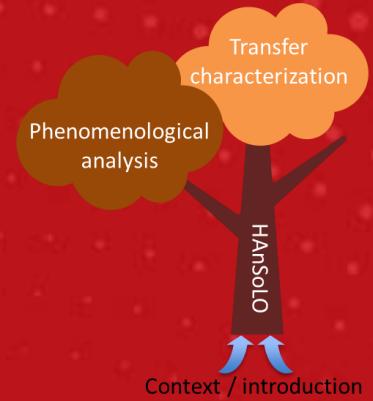
Stevens et al. (1991)



Shoulder in cavity shape = decrease of surface temperature
 Liquid film becomes turbulent



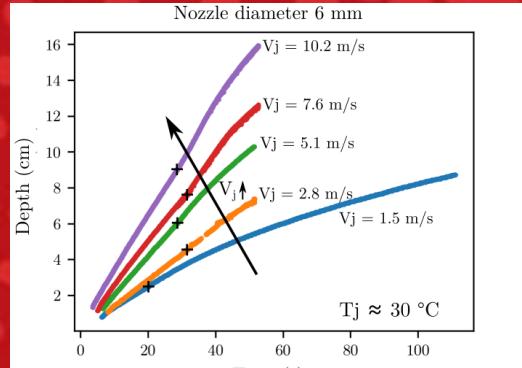
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Transfers characterization

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• Raw results example

• Film regime

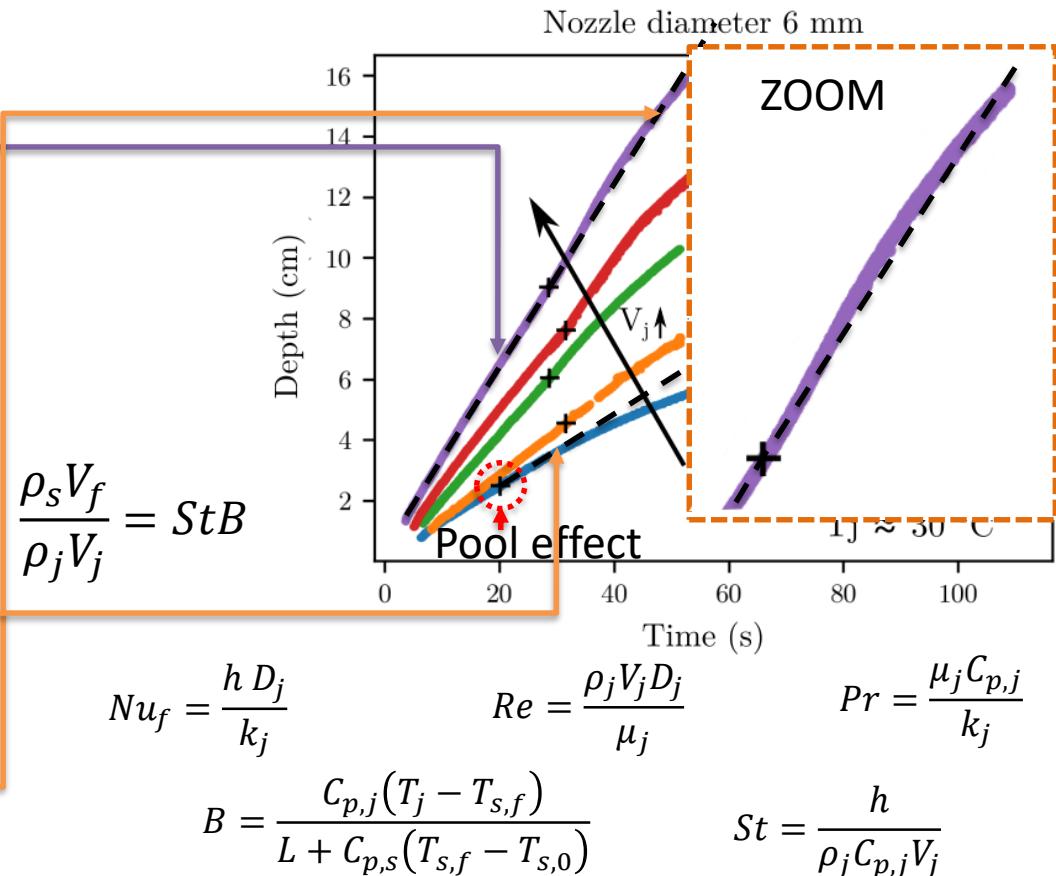
- Constant ablation velocity
- $V_f \uparrow$ with T_j & V_j
- $V_f \uparrow$ when $D_j \downarrow$
- h & Nu from jump condition + Newton's law + constant V_f

$$h = \frac{V_f \rho_s [L + C_{p,s}(T_{s,f} - T_{s,0})]}{(T_j - T_{s,f})}$$

$$Nu_f = K Re^n Pr^m$$

• Behavior change in pool effect

- Not linear
- Potential $V_f \uparrow$





• Transition to pool effect

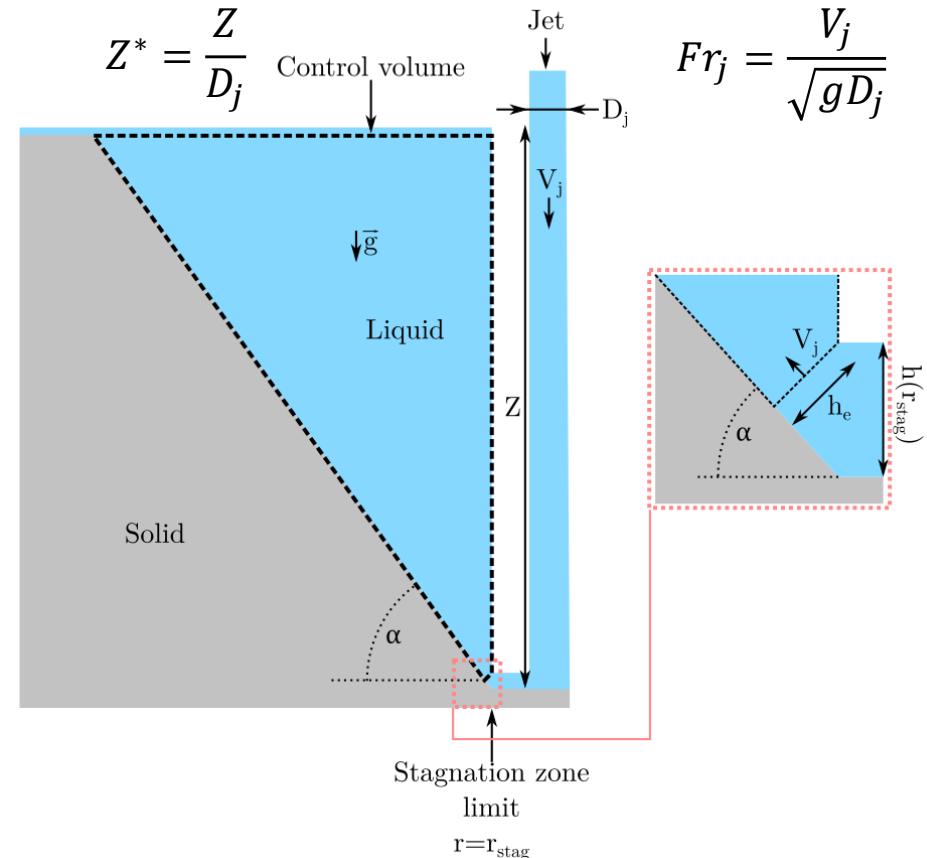
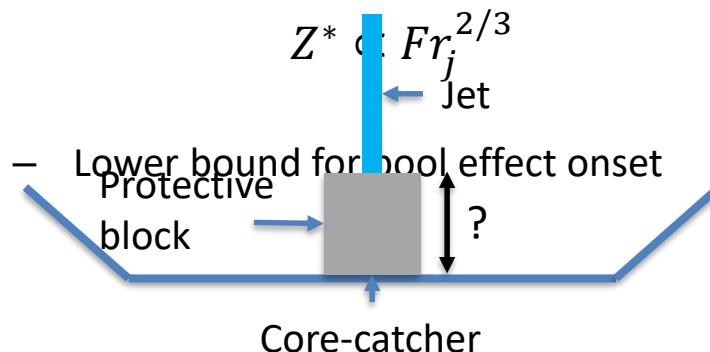
- 1st measurements
- Data well represented with Z^* & Fr_j
- Saito's et al. (1990) criterion does not work ($Z^* > 4$)
- Our model:

- Momentum balance on truncated cone

$$Z^* \propto Fr_j^{2/3}$$

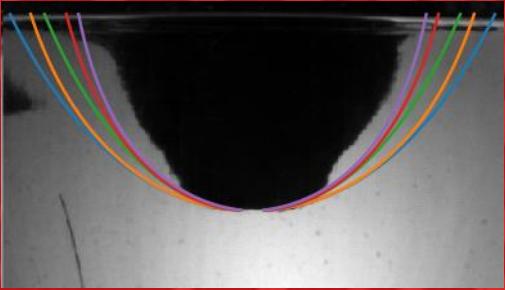
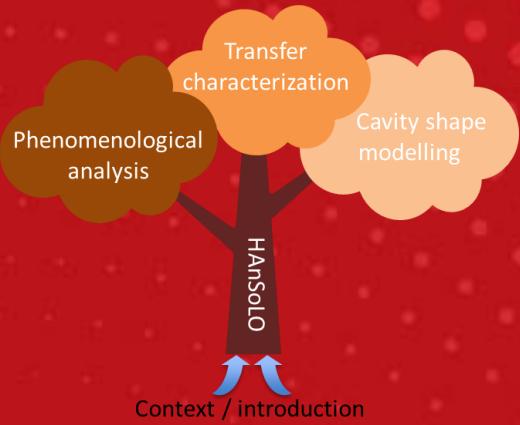
Jet

- Lower bound for pool effect onset





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Cavity shape modelling

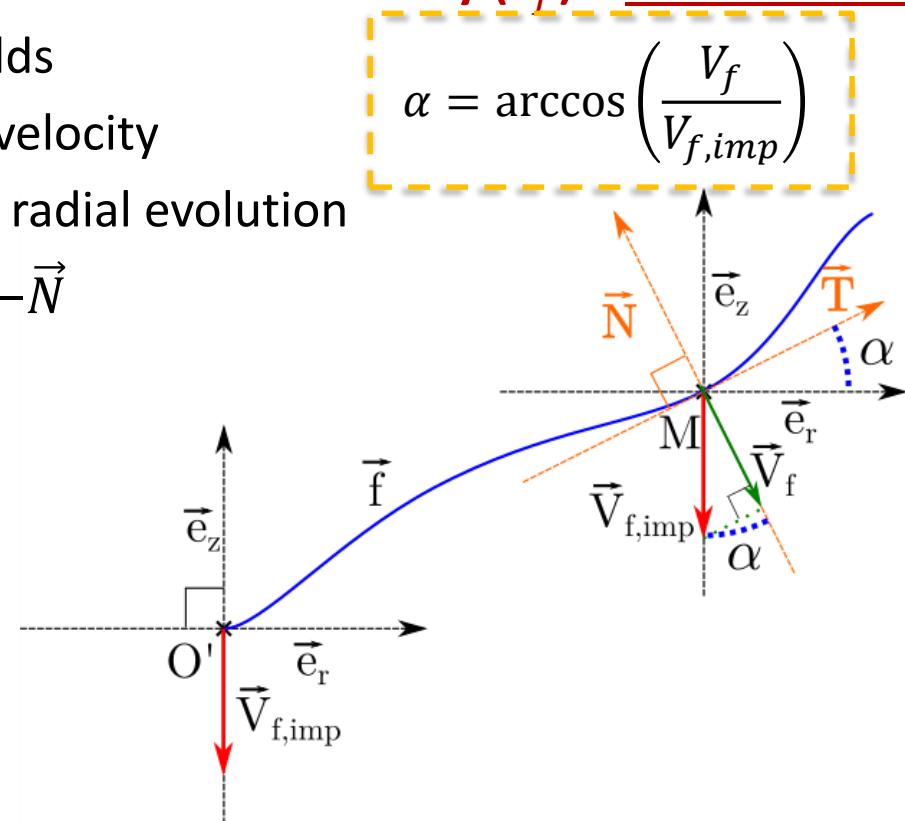
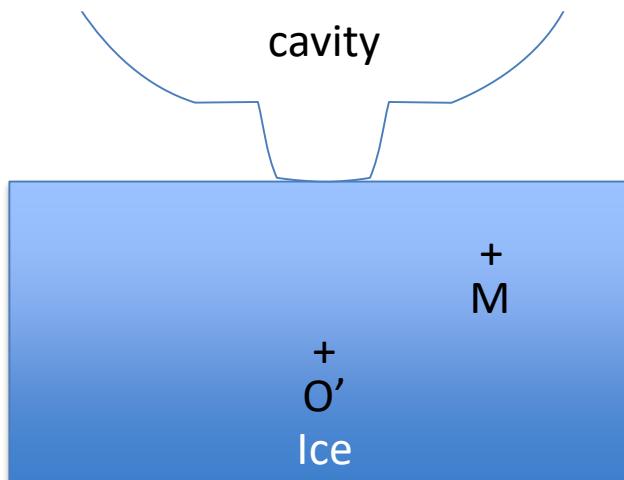
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• Link between cavity local angle (α) & local ablation velocity (V_f) – FILM REGIME

- Cavity maintains its shape : it unfolds
- Interface points move at constant velocity
- Cavity shape from melting velocity radial evolution
- Points displaced by melting along $-\vec{N}$





• Modelling: Boundary layer growth

• Constant film temperature

- IR measurements

- Only h evolve with r

$$V_f = \frac{h(T_{film} - T_{s,f})}{\rho_s [L + C_{p,s}(T_{s,f} - T_{s,0})]}$$



• Evolution from scaling law without melting Lienhard (2006):

- z evolution from α

$$Nu_{stag} = 0.745 Re^{1/2} Pr^{1/3}$$

$$Nu = 0.632 Re^{1/2} Pr^{1/3} \sqrt{\frac{D_j}{r}}$$

$$\frac{V_f}{V_{f,imp}} = \frac{h}{h_{stag}} = \frac{Nu}{Nu_{stag}}$$

$$\text{&} \quad z = \int_0^r \sqrt{\left(\frac{V_f}{V_{f,imp}}\right)^2 - 1} d\mathcal{L}$$

Cavity shape prevision



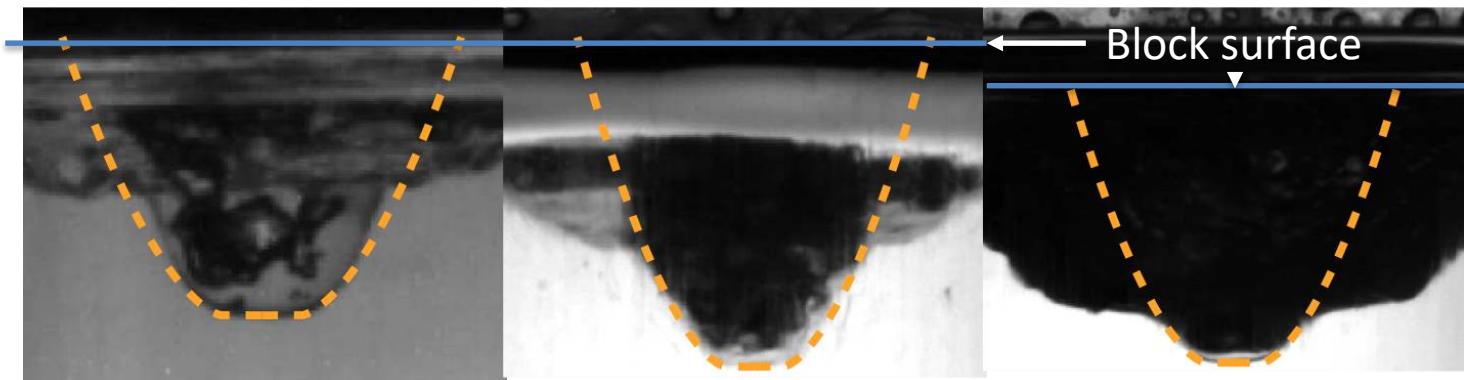
• Modelling: Boundary layer growth (comparisons)

- At pool effect transition: $T_j = 50^\circ C$, $D_N = 6.0 \text{ mm}$

$$\begin{aligned}V_j &= 1.2 \text{ m/s} \\Re &= 13\,000 \\Pr &= 3.70\end{aligned}$$

$$\begin{aligned}V_j &= 2.3 \text{ m/s} \\Re &= 24\,000 \\Pr &= 3.55\end{aligned}$$

$$\begin{aligned}V_j &= 5.1 \text{ m/s} \\Re &= 55\,000 \\Pr &= 3.55\end{aligned}$$



— —
Calculated shape



• Modelling: constant h - $T_{film} \downarrow$ modelling

Hypotheses

- T uniform along e
- Newton law at interface $\varphi = h(T - T_{s,f})$
- Steady state
- Perfect fluid
- No transfers in stagnation zone



Balances

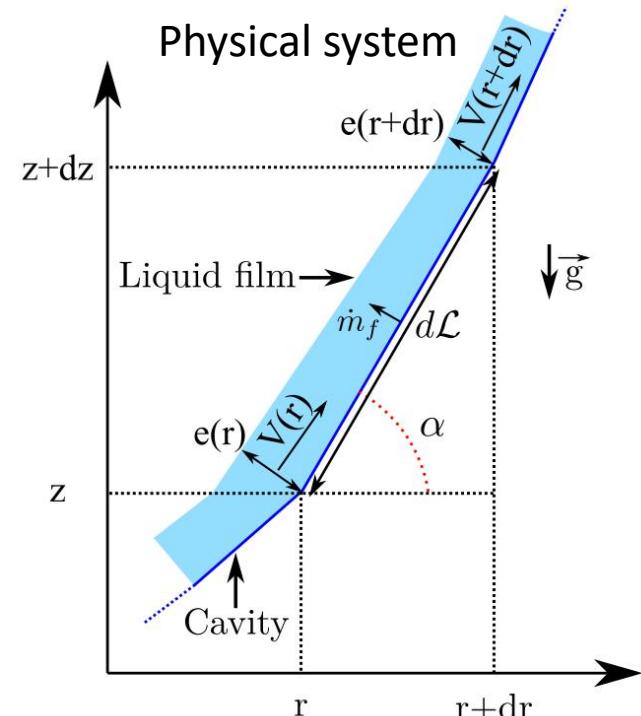


ODE system

- Mass
- Heat
- Momentum

3 Dimensionless numbers

$$B = \frac{C_{p,j}(T_j - T_{s,f})}{L + C_{p,s}(T_{s,f} - T_{s,0})} \quad St = \frac{h}{\rho_j C_{p,j} V_j} \quad Fr_j = \frac{V_j}{\sqrt{g D_j}}$$

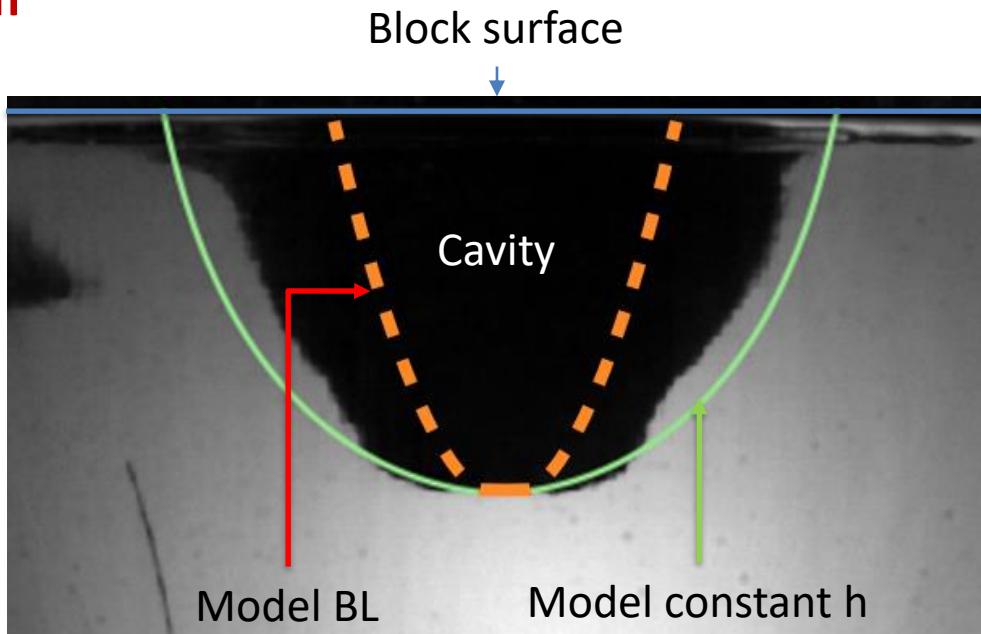




• Modelling: constant h - Comparison

- Model describes well global cavity shape with no laminar part
- Interesting estimates of cavity width

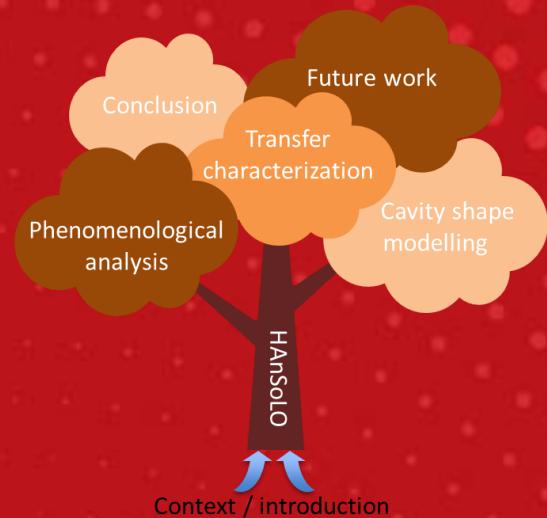
$$\left\{ \begin{array}{l} \frac{d}{dr^*}(r^* e^* V^*) = \frac{St B \theta}{\cos(\alpha)} \\ \frac{1}{r^*} \frac{d}{dr^*}(r^* e^* V^* \theta) = - \frac{St \theta}{\cos(\alpha)} \\ \frac{1}{r^*} \frac{d}{dr^*}(r^* e^* V^{*2}) \sin(\alpha) = - \frac{e^*}{Fr_j^2} \frac{1}{\cos(\alpha)} \\ \frac{dz^*}{dr^*} = \sqrt{\left(\frac{\theta}{\theta_0}\right)^2 - 1} \\ \alpha = \arccos\left(\frac{\theta}{\theta_0}\right) \end{array} \right.$$



$$\begin{array}{ll} T_j = 71 \text{ } ^\circ\text{C} & Re = 147 \text{ 000} \\ V_j = 10.1 \text{ m/s} & Pr = 2.54 \end{array}$$



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Conclusion and Current Work (Antoine Avrit)

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